

## Temperature dependence of dead time as determined by methane retention in open-tubular capillary gas chromatography

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### ABSTRACT

The temperature dependence of the retention time of methane was determined in open-tubular capillary gas chromatography with either nitrogen or hydrogen as the carrier gas. Linear regression according to the equation  $t_M = AT^B$  gave smaller standard deviations  $s$  for  $0.5 < B < 1$  than for  $B = 0.5$ . However, the equation  $t_M = Ae^{B/T}$  may fit almost equally well, the constant  $B$  being approximately independent of the column inlet pressure and of the nature of the carrier gas, which facilitates the determination of the theoretically important term  $\delta \ln k'/\delta(1/T)$ .

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### INTRODUCTION

Interest has been shown in the use of retention data for qualitative analysis ever since the advent of gas chromatography [1]. Several methods for the determination of the column dead time  $t_0$  are currently in use. The non-coincidence [2] of gas hold-up  $t_0$  and retention time  $t_M$  of a “non-interacting” compound such as air [3] or methane [4] has led to the determination of  $t_0$  by a variety of computational methods based on retention data of hydrocarbons, some of which involve a high degree of sophistication [5]. Whereas significant differences between  $t_M$  and  $t_0$  were found for packed columns showing methane adsorption [6,7], in gas-liquid chromatography with contemporary open-tubular capillary columns the methane retention time  $t_M$  in first order represents a fairly good approximation for  $t_0$ , provided that the column temperature is sufficiently high.

The temperature dependence of  $t_M$  is of interest in two fields of increasing importance, viz., the determination of thermodynamic parameters of solute-solvent interactions, particularly useful for chiral recognition studies [8–10], and the prediction of retention data under temperature programming [11–14]. Ettre [15] calculated temperature-dependent viscosity data for several carrier gases that were adopted by many workers [11–13,16–19]; however, only a few experimental data on the temperature dependence of  $t_M$  have been published [17,18]. This observation prompts us to report on some previous experimental investigations on this subject [20,21].

## EXPERIMENTAL

Gas chromatography was performed on a Carlo-Erba Model 2101 gas chromatograph with hydrogen or nitrogen as the carrier gas. The inlet pressure was kept at a set of constant values, as indicated; the splitting ratio was 1:50 and detection was effected with a flame ionization detector. Retention data were measured in triplicate, using a stop-watch. Linear regression calculations were performed according to the method of Deming [22], where each data point is represented by a three-dimensional Gaussian.

Non-deactivated borosilicate glass capillaries were coated with L-Chirasil-Val [23–25], as described [26]. A stainless-steel capillary column was coated with squalane, as described [27].

## RESULTS AND DISCUSSION

The capillary column may be compared with an Oswald viscometer; according to the Hagen–Poiseuille law let [28]

$$\eta = Cdt_M \quad (1)$$

where  $\eta$  = viscosity of the gas,  $d$  = density,  $C$  = constant for the given apparatus and  $t_M$  = gas hold-up. Neglecting the temperature dependence of  $C$  and  $d$ , it follows that

$$\frac{\delta t_M}{dT} \propto \frac{\delta \eta}{dT} \quad (2)$$

Hence the effect of the temperature dependence of dead time  $t_M$  is reduced to the well known problem of the temperature dependence of the viscosity  $\eta$ . Depending on the gas theory applied, several equations have been established [29]:

$$\eta \propto T^{0.5} \quad \text{ideal gas theory} \quad (3)$$

$$\eta \propto T^B \quad \text{for real gases} \quad (4)$$

where  $B$  varies from *ca.* 0.7 for hydrogen up to 1.0 for less ideal gases, and

$$\eta \propto \frac{T^{0.5}}{1 + C/T} \quad \text{Sutherland equation} \quad (5)$$

where  $C$  is a constant specific for the particular compound, and for  $C = 0$  eqn. 5 is reduced to eqn. 3. It is worth mentioning that the temperature dependence of the viscosity played a historical role in establishing a suitable model for gases, *e.g.*, Maxwell erroneously predicted an exponent  $B = 1.0$  in eqn. 4, assuming that the molecules would repel one another with a force inverse to the fifth power of the distance between them [29]. In contrast, the Sutherland eqn. 5, still in use today [28,30,31], is based on the assumption of a viscous action between successive gas layers.

TABLE I

TEMPERATURE DEPENDENCE OF THE RETENTION TIME OF METHANE ON SQUALANE WITH NITROGEN AS THE CARRIER GAS: COMPARISON OF EXPERIMENTAL AND CALCULATED DATA

Stainless-steel capillary column (100 m  $\times$  0.5 mm I.D.), inlet pressure 0.5 bar, splitting ratio 1:50, flame ionization detection; conditions (a) (see Table III).

| $T(^{\circ}\text{C})$ | $t_{\text{M}}(\text{s})$                     |                  |                  |                  |                  |
|-----------------------|--|------------------|------------------|------------------|------------------|
|                       | Experimental,<br>$t_{\text{M}_{\text{exp}}}$ | Calculated       |                  |                  |                  |
|                       |  | $t_{\text{M}_1}$ | $t_{\text{M}_2}$ | $t_{\text{M}_3}$ | $t_{\text{M}_4}$ |
| 50                    | 359.3  | 359.22           | 359.28           | 363.42           | 359.16           |
| 60                    | 367.6  | 367.68           | 367.68           | 369.00           | 367.74           |
| 70                    | 376.0  | 376.02           | 375.96           | 374.52           | 376.08           |
| 80                    | 384.2  | 384.18           | 384.18           | 379.92           | 384.06           |

In view of this discussion, we decided to determine  $t_{\text{M}}$  in a gas chromatographic apparatus at different temperatures. In a first experiment with nitrogen as the carrier gas (see Table I), experimental data  $t_{\text{M}_{\text{exp}}}$  were compared with  $t_{\text{M}_1}$ ,  $t_{\text{M}_2}$  and  $t_{\text{M}_4}$  obtained from different regression analyses:

$$t_{\text{M}_1} = AT^{0.5} + B \quad (6)$$

$$t_{\text{M}_2} = AT^B \quad (7)$$

$$t_{\text{M}_3} = AT^{0.5} \quad (8)$$

$$t_{\text{M}_4} = Ae^{B/T} \quad (9)$$

The one-parameter approximation, eqn. 8, can be ruled out by the fairly poor accordance with the experimental data for  $t_{\text{M}_{\text{exp}}}$  (see Table III). However, all two-parameter approximations examined, eqns. 6, 7 and 9, agree almost equally well with  $t_{\text{M}_{\text{exp}}}$ . On the basis of the standard deviations  $s$ , 0.06 s for both  $t_{\text{M}_1}$  and  $t_{\text{M}_2}$  and 0.15 s for  $t_{\text{M}_4}$ , a clear distinction between the three different mathematical models is not possible.

The good description of the temperature dependence of the retention time of methane by the exponential approximation, eqn. 9 deserves particular attention. To the best of our knowledge, there is no corresponding equation for the viscosity of gases in literature. However, with respect to the viscosity of liquids, Andrade [32] published eqn. 10, which was later modified by Vogel to give eqn. 11 [33]:

$$\eta \propto e^{B/T} \quad (10)$$

$$\eta \propto e^{B/(T-C)} \quad (11)$$

The underlying model is not too different from that discussed for the viscosity of gases in the way in which a given flux is assumed, and the interaction of one layer of molecules with an adjacent layer of molecules in the liquid is used to derive the viscosity behaviour. The high predictive power of Andrade's eqn. 10 was regarded as "quite astonishing" in view of the imperfections in all the considerations involved [33].

In contrast to this statement, one may argue that an open-minded comparison of the different theoretical approaches may eventually demonstrate the arbitrariness of the multi-parameter regression analyses applied, in view of the limited accuracy of the experimental data. In order to shed some more light on the problem, an extended series of measurements were performed with hydrogen as the carrier gas. For each of four different values of the inlet pressure fourteen data points were taken (see Table II). For  $p_i = 0.5$  bar hydrogen, the parabolic approximation  $t_{M_2}$ , calculated according to eqn. 7 as published previously [15], was compared with the exponential approximation  $t_{M_4}$ , as proposed in eqn. 9 (see Fig. 1). No significant advantage of the former approach over the latter can be seen from this plot.

A similar observation is made for the overall set of experimental data, as judged from the standard deviation  $s$ . The relative error in the temperature coefficient  $B$  is in the range 0.9–2.2% for eqn. 7 and 1.5–2.9% for eqn. 9. However, these deviations are far below the relative errors calculated for the mean values of  $B$ , i.e., 3.8% for eqn. 7 and 4.0% for eqn. 9. There is no apparent dependence of  $B$  on the inlet pressure  $p_i$ . As predicted by the classical kinetic gas theory,  $\eta$  of an ideal gas should be independent of its pressure  $p$  [29]. Moreover, there should be no apparent influence of the pressure drop  $\Delta p = (p_i - p_0)$  on the temperature coefficient  $B$ . This prediction is justified by the

TABLE II

TEMPERATURE AND INLET PRESSURE DEPENDENCE OF THE RETENTION TIME OF METHANE ON L-CHIRASIL-VAL WITH HYDROGEN AS THE CARRIER GAS (EXPERIMENTAL DATA)

Borosilicate glass capillary column (No. H14, 20 m  $\times$  0.25 mm I.D.), splitting ratio 1:50, flame ionization detection, conditions (b) (see Table III).

| $T(^{\circ}\text{C})$ | $t_{M_{exp}}$ (s) |                   |                   |                   |
|-----------------------|-------------------|-------------------|-------------------|-------------------|
|                       | 0.3               | 0.5               | 0.7               | 1.0 [bar]         |
| 50                    | 52.9 <sub>3</sub> | 32.0 <sub>7</sub> | 23.4 <sub>0</sub> | 17.3 <sub>0</sub> |
| 60                    | 53.8 <sub>0</sub> | 32.6 <sub>7</sub> | 23.8 <sub>3</sub> | 17.5 <sub>3</sub> |
| 70                    | 54.8 <sub>3</sub> | 33.2 <sub>0</sub> | 24.2 <sub>3</sub> | 17.9 <sub>3</sub> |
| 80                    | 55.3 <sub>7</sub> | 34.1 <sub>3</sub> | 24.5 <sub>3</sub> | 18.2 <sub>0</sub> |
| 90                    | 56.4 <sub>8</sub> | 34.6 <sub>3</sub> | 25.0 <sub>7</sub> | 18.6 <sub>0</sub> |
| 100                   | 57.5 <sub>3</sub> | 35.1 <sub>7</sub> | 25.4 <sub>5</sub> | 18.8 <sub>8</sub> |
| 110                   | 58.8 <sub>3</sub> | 35.7 <sub>3</sub> | 25.9 <sub>5</sub> | 19.2 <sub>7</sub> |
| 120                   | 59.9 <sub>3</sub> | 36.3 <sub>0</sub> | 26.5 <sub>0</sub> | 19.5 <sub>5</sub> |
| 130                   | 61.2 <sub>0</sub> | 37.0 <sub>7</sub> | 27.2 <sub>3</sub> | 19.9 <sub>7</sub> |
| 140                   | 62.1 <sub>3</sub> | 37.9 <sub>0</sub> | 27.7 <sub>7</sub> | 20.2 <sub>3</sub> |
| 150                   | 62.9 <sub>3</sub> | 38.1 <sub>7</sub> | 28.1 <sub>0</sub> | 20.4 <sub>7</sub> |
| 160                   | 64.1 <sub>0</sub> | 38.5 <sub>5</sub> | 28.3 <sub>0</sub> | 20.7 <sub>7</sub> |
| 170                   | 65.0 <sub>3</sub> | 39.1 <sub>3</sub> | 28.8 <sub>7</sub> | 21.0 <sub>7</sub> |
| 180                   | 66.0 <sub>7</sub> | 39.8 <sub>3</sub> | 29.2 <sub>7</sub> | 21.4 <sub>0</sub> |

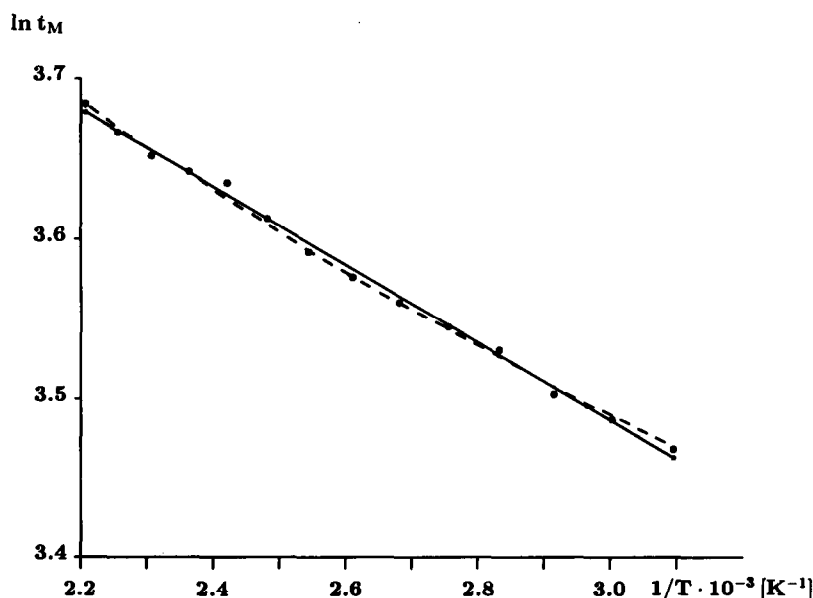


Fig. 1. Comparison of the two approaches,  $t_{M_2}$  and  $t_{M_4}$ , with the experimental data; conditions (b) (Table III), 0.5 bar  $H_2$ . Solid line,  $t_{M_2} = 68.2e^{-246/T}$  s; dashed line,  $t_{M_2} = 0.80T^{0.64}$  s.

TABLE III  
LINEAR REGRESSION ANALYSES

| Chromatographic conditions <sup>a</sup>                              | Theoretical model        | Regression equation for $t_M$ (s)                              | $s$ (s) <sup>b</sup> |
|--|--------------------------|--|----------------------|
| (a) 0.5 bar $N_2$  | $t_{M_1} = AT^{0.5} + B$ | $t_{M_1} = (30.55 \pm 0.09)\sqrt{T} - (190.0 \pm 1.6)$         | 0.06                 |
|  | $t_{M_2} = AT^B$         | $\ln t_{M_2} = (0.756 \pm 0.002) \ln T + (1.517 \pm 0.012)$    | 0.06                 |
|  | $t_{M_3} = AT^{0.5}$     | $t_{M_3} = (20.22 \pm 0.20)\sqrt{T}$                           | 4.46                 |
|  | $t_{M_4} = Ae^{B/T}$     | $\ln t_{M_4} = -(255.2 \pm 2.2) \cdot 1/T + (6.673 \pm 0.007)$ | 0.15                 |
| (b) 0.3 bar $H_2$<br>0.5 bar $H_2$<br>0.7 bar $H_2$<br>1.0 bar $H_2$ | $t_{M_2} = AT^B$         | $\ln t_{M_2} = (0.675 \pm 0.011) \ln T + (0.059 \pm 0.063)$    | 0.24                 |
|  |                          | $\ln t_{M_2} = (0.640 \pm 0.010) \ln T - (0.227 \pm 0.061)$    | 0.14                 |
|  |                          | $\ln t_{M_2} = (0.688 \pm 0.015) \ln T - (0.830 \pm 0.091)$    | 0.15                 |
|  |                          | $\ln t_{M_2} = (0.638 \pm 0.006) \ln T - (0.841 \pm 0.038)$    | 0.05                 |
|  | Mean value for $B$       | $(0.660 \pm 0.025)$  |                      |
| (b) 0.3 bar $H_2$<br>0.5 bar $H_2$<br>0.7 bar $H_2$<br>1.0 bar $H_2$ | $t_{M_4} = Ae^{B/T}$     | $\ln t_{M_4} = -(261.4 \pm 6.5) \cdot 1/T + (4.762 \pm 0.017)$ | 0.31                 |
|  |                          | $\ln t_{M_4} = -(246.0 \pm 4.3) \cdot 1/T + (4.223 \pm 0.011)$ | 0.04                 |
|  |                          | $\ln t_{M_4} = -(264.6 \pm 7.6) \cdot 1/T + (3.957 \pm 0.020)$ | 0.05                 |
|  |                          | $\ln t_{M_4} = -(245.2 \pm 3.6) \cdot 1/T + (3.600 \pm 0.009)$ | 0.05                 |
|  | Mean value for $B$       | $-(255.3 \pm 10.2)$  |                      |

<sup>a</sup> As indicated in Tables I and II.

<sup>b</sup> Standard deviation  $s = \sqrt{\frac{1}{n-2} \sum_{i=1}^n (t_{M_{exp}} - t_{M_j})^2}$

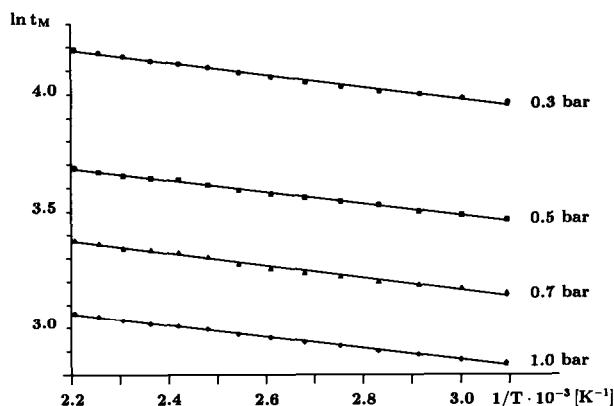


Fig. 2. Comparison of experimental and theoretical data  $t_{M_2}$ ; conditions (b) (Table III), 0.3, 0.5, 0.7 and 1.0 bar  $H_2$ .

data in Table III. The entire data set of conditions (b) is depicted in Fig. 2, together with the regression analyses according to the exponential eqn. 9.

Hence this approach proves to be equally valid for gases, and not only for liquids as originally assumed. Intriguingly, both carrier gases, hydrogen and nitrogen, show a similar temperature coefficient, *i.e.*,  $B = -255$ . Whether this coincidence is real or an artefact remains open to further investigations. In the equation  $t_M = AT^B$ , a significant difference in  $B$  is found for the different gases, *i.e.*,  $B = 0.756$  for nitrogen and  $B = 0.660$  for hydrogen. Both values differ only slightly from those used by Ettre [15] for the temperature coefficient of the viscosity  $\eta \propto T^B$  with  $B = 0.725$  for nitrogen and  $B = 0.680 - 0.695$  for hydrogen.

## CONCLUSIONS

The temperature dependence of the retention time of methane is in good accord with both eqns. 7 and 9. As judged by the standard deviations  $s$ , both approaches are of striking resemblance. The latter, although following a viscosity theory originally established by Andrade [32] for liquids, has the advantage of fitting into the determination of thermodynamic parameters of solute-solvent interaction by gas chromatography. We found it convenient simply to determine the temperature dependence of the net retention  $t'_R$  at a constant inlet pressure [9], which yields the theoretically important term

$$\frac{\delta \ln k'}{\delta(1/T)} = \frac{\delta \ln t'_R}{\delta(1/T)} - B \quad (12)$$

where the constant  $B$  is approximately independent of the column inlet pressure and of the nature of the carrier gas. Beyond these applications in the realm of chromatography, the exponential equation  $\eta \propto e^{B/T}$  proposed for the temperature dependence of the viscosity of gases may be of general interest with regard to model building in kinetic gas theory [29].

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